







Optimal control in finite strain elasticity and optimal design in electromechanics

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Online Seminar Control in Times of Crisis

Control in finite elasticity

- Finite Strain Elasticity
- The optimal control problem
- Numerical examples

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- Introduction
- Background
- The Design problem
- Numerical examples

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Design in Electro-Mechanics

Finite Strain Elasticity

Large strain kinematics: F, H, J



$$F = \nabla_0 \phi;$$
 $J = \det F;$ $H = JF^{-T}$

Finite Strain Elasticity

Finite Strain Elasticity: the BVP

· Tensor cross product [de Boer, 1982; Bonet, Gil and Ortigosa, 2015]:

$$(\boldsymbol{A} \times \boldsymbol{B})_{ij} = \mathcal{E}_{ikl} \mathcal{E}_{jmn} A_{km} B_{lm}$$

The area and volume maps are redefined as:

$$H = \frac{1}{2} F \times F; \qquad J = \frac{1}{6} (F \times F) : F$$

Strong form (PDE) in the Lagrangian formalism

$$\begin{aligned} \nabla_0 \times F &= 0; & \text{in } \Omega_0 & \Rightarrow F &= \nabla_0 \phi \\ \nabla_0 \cdot P(F) &+ b_0 &= 0; & \text{in } \Omega_0 & P &= \sigma H \\ \phi &= \phi^*; & \text{on } \partial\Omega_{0,D} \\ PN &= t_0; & \text{on } \partial\Omega_{0,N} \end{aligned}$$





Constitutive relation through nonlinear stored energy functional e(F):

$$P(F) = rac{\partial e(F)}{\partial F}$$

How to define e(F) being mathematically and physically admissible?

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Finite Strain Elasticity

Finite Strain Elasticity: polyconvexity

- · Polyconvexity conditon [Ball, 1976; Schröder and Neff, 2004...]
- The internal strain energy functional *e* is expressed as a **CONVEX** multi-variable functional *W* in terms of **fibre**, **area** and **volume** maps:

$$e\left(F
ight) = \mathbb{W}\left(\mathcal{V}
ight); \qquad \mathcal{V} = \{F, H(F), J(F)\}$$

$$\mathbb{W}\left(\lambda\boldsymbol{\mathcal{V}}_{1}+\left(1-\lambda\right)\boldsymbol{\mathcal{V}}_{2}\right)\leq\lambda\mathbb{W}\left(\boldsymbol{\mathcal{V}}_{1}\right)+\left(1-\lambda\right)\mathbb{W}\left(\boldsymbol{\mathcal{V}}_{2}\right)\quad\lambda\in\left[0,1\right],\forall\boldsymbol{\mathcal{V}}_{1},\boldsymbol{\mathcal{V}}_{2}\in\boldsymbol{\mathcal{V}}$$

Polyconvex constitutive models:

$$\begin{aligned} \mathsf{NH} \ \ &\mathbb{W}\left(F,J\right) = \frac{\mu}{2} ||F||^{2} + f\left(J\right); & f\left(J\right) = -\mu \ln J + \frac{\lambda}{2} \left(J - 1\right)^{2} \\ \mathsf{MR} \ \ &\mathbb{W}\left(F,H,J\right) = \frac{\alpha}{2} ||F||^{2} + \frac{\beta}{2} ||H||^{2} + g\left(J\right); & g\left(J\right) = -\left(\alpha + 2\beta\right) \ln J + \frac{\lambda}{2} \left(J - 1\right)^{2} \\ \mathsf{Gent} \ &\mathbb{W}\left(F,J\right) = -\frac{\mu J_{m}}{2} \ln \left(1 - \frac{||F||^{2} - 3}{J_{m}}\right) + f\left(J\right); \end{aligned}$$

Finite Strain Elasticity

Finite Strain Elasticity: the weak form

The PDE and boundary conditions

$$\nabla_0 \times F = \mathbf{0} \text{ in } \Omega_0 \quad \Rightarrow \quad F = \nabla_0 \phi$$
$$\nabla_0 \cdot P + \mathbf{b}_0 = \mathbf{0} \text{ in } \Omega_0; \qquad P = \boldsymbol{\sigma} H$$
$$\phi = \overline{\phi} \text{ on } \partial \Omega_{0,D}$$
$$PN = T \text{ on } \partial \Omega_{0,N}$$



For $\phi \in H^1(\Omega_0)$, multiplication by test function $\delta \phi \in H^1_0(\Omega_0)$ leads to the weak form

$$\mathcal{W}(\phi) = \underbrace{\int_{\Omega_0} P(F) : \nabla_0 \delta \phi \, dV}_{\text{Elastic work}} - \underbrace{\left(\int_{\Omega_0} f_0 \cdot \delta \phi \, dV + \int_{\partial \Omega_{0N}} t_0 \cdot \delta \phi \, dA\right)}_{\text{External work}} = 0$$

Stationary condition of energy functional:

$$\Pi(\phi) = \inf_{\phi} \left\{ \underbrace{\int_{\Omega_0} \mathbb{W}(F, H, J) \, dV}_{\text{Elastic energy}} - \underbrace{\left(\int_{\Omega_0} f_0 \cdot \phi \, dV + \int_{\partial \Omega_{0N}} t_0 \cdot \phi \, dA \right)}_{\text{external energy}} \right\}$$

Control finite elasticity

Finite Strain Elasticity

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Control finite elasticity ○○○○●○○○○○ Design in Electro-Mechanics

Finite Strain Elasticity

Finite Strain Elasticity: existence of minimisers

Energy functional to be minimised:

$$\Pi(\boldsymbol{\phi}) = \inf_{\boldsymbol{\phi}} \left\{ \underbrace{\int_{\Omega_0} \mathbb{W}(\boldsymbol{F}, \boldsymbol{H}, \boldsymbol{J}) \, d\boldsymbol{V}}_{\text{Elastic energy}} - \underbrace{\left(\int_{\Omega_0} f_0 \cdot \boldsymbol{\phi} \, d\boldsymbol{V} + \int_{\partial \Omega_{0N}} t_0 \cdot \boldsymbol{\phi} \, d\boldsymbol{A}\right)}_{\text{external energy}} \right\}$$

[1] Polyconvexity: for $\mathcal{V} = \{A, B, c\}, A, B \in \mathbb{R}^{3 \times 3}, c \in \mathbb{R}^+$

$$\mathbb{W}\left(\lambda \boldsymbol{\mathcal{V}}_{1}+\left(1-\lambda\right)\boldsymbol{\mathcal{V}}_{2}\right) \leq \lambda \mathbb{W}\left(\boldsymbol{\mathcal{V}}_{1}\right)+\left(1-\lambda\right)\mathbb{W}\left(\boldsymbol{\mathcal{V}}_{2}\right) \quad \lambda \in \left[0,1\right], \forall \boldsymbol{\mathcal{V}}_{1}, \boldsymbol{\mathcal{V}}_{2} \in \boldsymbol{\mathcal{V}}$$

[2] Coerciveness: for $p \ge 2$, $\frac{1}{p} + \frac{1}{q} \le 1$, $r \ge 1$

$$\mathbb{W}\left(oldsymbol{\mathcal{V}}
ight) \geq lpha\left(||oldsymbol{F}||^{p}+||oldsymbol{H}||^{q}+oldsymbol{J}^{r}
ight) +eta \qquad lpha>0$$

[3] Limit behaviour:

$$\lim_{J\to 0^+} \mathbb{W}\left(\mathbf{F}, \mathbf{H}, \mathbf{J}\right) = +\infty$$

The optimal control problem

Biomimetic actuation

· Biomimetic-type actuation

$$\mathbb{W}_{\mathsf{fiber}}(\mathbf{F}, m) = \mathbb{W}(\mathbf{F}) + \mathbb{W}_{\mathsf{m}}(\mathbf{F}, m)$$

$$\begin{split} \mathbb{W}_{\mathsf{m}}(\boldsymbol{F}, m) &= -\frac{m}{2} ||\boldsymbol{F}\boldsymbol{a}||^2 \\ \boldsymbol{a} : \boldsymbol{a}(\boldsymbol{X}) : \ \Omega_0 \to \mathbb{R}^3 \\ m : m(\boldsymbol{X}) : \ \Omega_0 \to \mathbb{R} \end{split}$$

$$\boldsymbol{P}(\boldsymbol{F},m) = \partial_{\boldsymbol{F}} \mathbb{W}(\boldsymbol{F}) - m\boldsymbol{F}\boldsymbol{a} \otimes \boldsymbol{a}$$

• For Mooney-Rivlin model, polyconvexity and coerciveness guaranteed if:

$$\mathbb{W}(\boldsymbol{F}, \boldsymbol{H}, \boldsymbol{J}) = \frac{\alpha}{2} ||\boldsymbol{F}||^{2} + \frac{\beta}{2} ||\boldsymbol{H}||^{2} + g(\boldsymbol{J})$$
$$\sup_{\boldsymbol{X} \in \Omega_{0}} m(\boldsymbol{X}) < \alpha$$

· Stationary condition of energy functional:

$$\Pi(\phi, m) = \inf_{\phi} \left\{ \underbrace{\int_{\Omega_0} \mathbb{W}_{\text{fiber}}(F, H, J, m) \, dV}_{\text{Elastic energy}} - \underbrace{\left(\int_{\Omega_0} f_0 \cdot \phi \, dV + \int_{\partial \Omega_{0N}} t_0 \cdot \phi \, dA\right)}_{\text{external energy}} \right\}$$





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The optimal control problem I

The Optimal Control problem

 $\begin{cases} \mbox{Minimize in } m: & \mathcal{J}(m) \\ \mbox{subject to} & m \in L^2(\Omega_0) \\ & m(X) \leq C \mbox{ a.e. } X \in \Omega_0 \\ & \phi \mbox{ is a minimizer of } \Pi(\phi,m) \end{cases}$



A Hausdorff distance-based objective function

 $\mathcal{J}(m) = \rho_H\left(\Omega(m), \Omega_d\right)$

$$\rho_{H}\left(\Omega(m),\Omega_{d}\right) := \max\left\{\sup_{x\in\Omega_{d}}d_{\Omega(m)}\left(x\right),\sup_{x\in\Omega(m)}d_{\Omega_{d}}\left(x\right)\right\}$$
$$d_{\Omega}\left(x\right) := \inf_{y\in\Omega}\left\|y-x\right\|$$

Regularised Hausdorff functional

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$$d_{\Omega} \left(\boldsymbol{x} \right) := \inf_{\boldsymbol{y} \in \Omega} \| \boldsymbol{y} - \boldsymbol{x} \|$$

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$$d_{\Omega} \left(\boldsymbol{x} \right) := \inf_{\boldsymbol{y} \in \Omega} \| \boldsymbol{y} - \boldsymbol{x} \|$$

Regularised Hausdorff functional

$$\mathcal{J}(m) = \tilde{\rho}_H\left(\Omega(m), \Omega_d\right)$$

The optimal control problem

The optimal control problem II

The Optimal Control problem

$$\left\{ \begin{array}{ll} \text{Minimize in } m: \quad \mathcal{J}(m) \\ \text{subject to} \qquad m \in L^2(\Omega_0) \\ \qquad m(\textbf{X}) \leq C \text{ a.e. } \textbf{X} \in \Omega_0 \\ \qquad \phi \text{ is a minimizer of } \Pi(\phi,m) \end{array} \right.$$



$$\mathcal{L}\left(\overline{\phi},\overline{p},\overline{m}\right) = \mathcal{J}\left(\overline{\phi},\overline{m}\right) - \int_{\Omega_{0}} P(\boldsymbol{\nabla}_{0}\overline{\phi},\overline{m}) : \boldsymbol{\nabla}_{0}\overline{p} \, dX$$
$$\mathcal{J}\left(\overline{\phi},\overline{m}\right) = \tilde{\rho}_{H}\left(\overline{\phi}(\Omega_{0}),\Omega_{d}\right) + \frac{M}{2} \int_{\Omega_{0}} \overline{m}^{2}\left(X\right) \, dX + \frac{\varepsilon}{2} \int_{\Omega_{0}} \left|\boldsymbol{\nabla}_{0}\overline{m}\left(X\right)\right|^{2} dX$$

$$\begin{split} \frac{\partial \mathcal{L}}{\partial \overline{p}} (\overline{\phi}, \overline{p}, \overline{m})(\mathbf{v}) &= \frac{\partial \Pi}{\partial \overline{\phi}} (\overline{\phi}, \overline{m})(\mathbf{v}) = \int_{\Omega_0} \mathbf{P}(\nabla_0 \overline{\Phi}, \overline{m}) : \nabla_0 \mathbf{v} \, dX = 0 \quad \forall \mathbf{v} \in H_0^1(\Omega_0) \\ \frac{\partial \mathcal{L}}{\partial \overline{\Phi}} (\overline{\Phi}, \overline{p}, \overline{m})(\mathbf{v}) &= \frac{\partial \tilde{\rho}_H}{\partial \overline{\Phi}} (\overline{\Phi}) (\mathbf{v}) - \int_{\Omega_0} \nabla_0 \overline{p} : \mathcal{C}(\nabla_0 \overline{\Phi}, \overline{m}) : \nabla_0 \mathbf{v} \, dX = 0 \quad \forall \mathbf{v} \in H_0^1(\Omega_0) \\ \mathcal{C}_{iijJ} &= \left(\nabla_{FF}^2 W(F, H(F), J(F)) \right)_{iijJ} - \delta_{ij} ma_I a_J \\ \frac{\partial \mathcal{L}}{\partial \overline{m}} (\overline{\Phi}, \overline{p}, \overline{m})(\hat{m}) = \int_{\Omega_0} \left[\hat{m} \left((Fa \otimes a) : \nabla_0 p \right) + M \, \overline{m} \, \hat{m} + \varepsilon \, \nabla_0 \overline{m} \cdot \nabla_0 \hat{m} \right] dX \end{split}$$

Numerical examples

Numerical example I



R. Ortigosa, J. Martínez-Frutos, C. Mora-Corral, P. Pedregal and F. Periago. Optimal control of soft materials using a Hausdorff distance functional, SICON, accepted.

Design in Electro-Mechanics

Numerical examples

Numerical example II



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Introduction

Finite Strain Electromechanics: motivation



Physical principle: Coulomb forces



Giant electrically induced deformations in EAP membrane [Suo 2013 (Harvard Lab), Max Plank Press 2016]



Control finite elasticity

Background

Finite Strain Electromechanics: the BVP

• Material strong form of **electrostatics** (e) in terms of E_0, D_0 :

$$\begin{aligned} \boldsymbol{\nabla}_0 \times \boldsymbol{E}_0 &= \boldsymbol{0} \text{ in } \Omega_0 \quad \Rightarrow \quad \boldsymbol{E}_0 &= -\boldsymbol{\nabla}_0 \varphi = \boldsymbol{F}^T \boldsymbol{E} \\ \boldsymbol{\nabla}_0 \cdot \boldsymbol{D}_0 &- \rho_0 &= 0 \text{ in } \Omega_0; \qquad \boldsymbol{D}_0 &= \boldsymbol{H}^T \boldsymbol{D} \\ \varphi &= \overline{\varphi} \text{ on } \partial \Omega_{0,De} \\ \boldsymbol{D}_0 \cdot \boldsymbol{N} &= -\omega_0 \text{ on } \partial \Omega_{0,Ne} \end{aligned}$$



Material strong form for the mechanics (m) in terms of P, F

$$\begin{aligned} \nabla_0 \times F &= \mathbf{0} \text{ in } \Omega_0 \implies F = \nabla_0 \phi \\ \nabla_0 \cdot P + b_0 &= \mathbf{0} \text{ in } \Omega_0; \qquad P = \sigma H \\ \phi &= \overline{\phi} \text{ on } \partial \Omega_{0,Dm} \\ PN &= T \text{ on } \partial \Omega_{0,Nm} \end{aligned}$$

With a coupling constitutive law [Ogden & Dorlmann, Bustemente, McMeeking & Landis, Suo & Zheo]

$$e\left(F,D_{0}\right)=c_{m}\left(F\right)+c_{c}\left(F,D_{0}\right)$$

Control finite elasticity

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$$\begin{aligned} \boldsymbol{\nabla}_0 \times \boldsymbol{F} &= \boldsymbol{0} \text{ in } \Omega_0 \quad \Rightarrow \quad \boldsymbol{F} &= \boldsymbol{\nabla}_0 \boldsymbol{\phi} \\ \boldsymbol{\nabla}_0 \cdot \boldsymbol{P} &+ \boldsymbol{b}_0 &= \boldsymbol{0} \text{ in } \Omega_0; \qquad \boldsymbol{P} &= \boldsymbol{\sigma} \boldsymbol{H} \\ \boldsymbol{\phi} &= \overline{\boldsymbol{\phi}} \text{ on } \partial \Omega_{0,Dm} \\ \boldsymbol{PN} &= \boldsymbol{T} \text{ on } \partial \Omega_{0,Nm} \end{aligned}$$



With a coupling constitutive law [Ogden & Dorfmann, Bustamante, McMeeking & Landis, Suo & Zhao]

$$e(F, D_0) = e_m(F) + e_e(F, D_0)$$

Constitutive relations:

$$P = rac{\partial e(F, D_0)}{\partial F}$$
 $E_0 = rac{\partial e(F, D_0)}{\partial D_0}$

Background

Finite Strain Electromechanics: the BVP

Material strong form of **electrostatics** (e) in terms of E_0, D_0 :

$$\begin{aligned} \boldsymbol{\nabla}_0 \times \boldsymbol{E}_0 &= \boldsymbol{0} \text{ in } \Omega_0 \quad \Rightarrow \quad \boldsymbol{E}_0 &= -\boldsymbol{\nabla}_0 \varphi = \boldsymbol{F}^T \boldsymbol{E} \\ \boldsymbol{\nabla}_0 \cdot \boldsymbol{D}_0 &- \rho_0 &= 0 \text{ in } \Omega_0; \qquad \boldsymbol{D}_0 &= \boldsymbol{H}^T \boldsymbol{D} \\ \varphi &= \overline{\varphi} \text{ on } \partial \Omega_{0,De} \\ \boldsymbol{D}_0 \cdot \boldsymbol{N} &= -\omega_0 \text{ on } \partial \Omega_{0,Ne} \end{aligned}$$

Material strong form for the **mechanics** (*m*) in terms of **P**, **F**:

$$\begin{aligned} \boldsymbol{\nabla}_0 \times \boldsymbol{F} &= \boldsymbol{0} \text{ in } \Omega_0 \quad \Rightarrow \quad \boldsymbol{F} &= \boldsymbol{\nabla}_0 \boldsymbol{\phi} \\ \boldsymbol{\nabla}_0 \cdot \boldsymbol{P} &+ \boldsymbol{b}_0 &= \boldsymbol{0} \text{ in } \Omega_0; \qquad \boldsymbol{P} &= \boldsymbol{\sigma} \boldsymbol{H} \\ \boldsymbol{\phi} &= \overline{\boldsymbol{\phi}} \text{ on } \partial \Omega_{0,Dm} \\ \boldsymbol{PN} &= \boldsymbol{T} \text{ on } \partial \Omega_{0,Nm} \end{aligned}$$

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$$e(\mathbf{F},\mathbf{D}_0) = e_m(\mathbf{F}) + \frac{\mathbf{e}_{\mathbf{e}}}{\mathbf{e}_{\mathbf{e}}}(\mathbf{F},\mathbf{D}_0)$$

Constitutive relations:

$$P = rac{\partial e\left(F, D_0
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$$e(F, D_0) = e_m(F) + \frac{e_e}{e}(F, D_0)$$

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$$\boldsymbol{P} = \frac{\partial e\left(\boldsymbol{F}, \boldsymbol{D}_{0}\right)}{\partial \boldsymbol{F}} \qquad \boldsymbol{E}_{0} = \frac{\partial e\left(\boldsymbol{F}, \boldsymbol{D}_{0}\right)}{\partial \boldsymbol{D}_{0}}$$

Control finite elasticity

Background

Finite Strain Electromechanics: polyconvexity

· Classical approach [Ogden & Dorfmann, Bustamante, McMeeking & Landis, Suo & Zhao, Vu & Steinman]

$$e(\mathbf{F}, \mathbf{D}_0) = e_m(\mathbf{F}) + e_e(\mathbf{F}, \mathbf{D}_0)$$

For **polyconvex** reversible electromechanics:

$$e(\mathbf{F}, \mathbf{D}_0) = \mathbb{W}(\mathbf{\mathcal{V}}); \qquad \mathbf{\mathcal{V}} = \{\mathbf{F}, \mathbf{H}, \mathbf{J}, \mathbf{D}_0, d\}; \qquad d = \mathbf{F} \mathbf{D}_0$$

 $\mathbb{W}\left(\lambda \boldsymbol{\mathcal{V}}_{1}+\left(1-\lambda\right) \boldsymbol{\mathcal{V}}_{2}\right) \leq \lambda \mathbb{W}\left(\boldsymbol{\mathcal{V}}_{1}\right)+\left(1-\lambda\right) \mathbb{W}\left(\boldsymbol{\mathcal{V}}_{2}\right) \quad \lambda \in \left[0,1\right], \forall \boldsymbol{\mathcal{V}}_{1}, \boldsymbol{\mathcal{V}}_{2} \in \boldsymbol{\mathcal{V}}$

A possible polyconvex electrically enhanced model is:

$$\mathbb{W}(F, H, J, D_0, d) = \underbrace{\mu_1 ||F||^2 + \mu_2 ||H||^2 + f(J)}_{\text{Mooney-Rivlin model}} + \underbrace{\frac{1}{2\varepsilon_2} ||D_0||^2 + \frac{1}{2\varepsilon_1} ||d||^2}_{\text{Electro modernical term}}$$

Constitutive relations:

$$P = \frac{\partial e(F, D_0)}{\partial F} \qquad P = \frac{\partial \mathbb{W}}{\partial F} + \frac{\partial \mathbb{W}}{\partial H} \times F + \frac{\partial \mathbb{W}}{\partial J} H + \frac{\partial \mathbb{W}}{\partial d} \otimes D_0$$
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A. J. Gil and R. Ortigosa. A new framework for large strain electromechanics based on convex multi-variable strain energies: Variational formulation and material characterisation, CMAME 302 (2016) 293-328.

Control finite elasticity

Background

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Control finite elasticity

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A. J. Gil and R. Ortigosa. A new framework for large strain electromechanics based on convex multi-variable strain energies: Variational formulation and material characterisation, CMAME 302 (2016) 293-328.

Control finite elasticity

Background

Topology optimisation of Electro Active material



R. Ortigosa, J. Martínez-Frutos, D. Ruíz, A. Donoso and J. C. Bellido. Density-based Topology Optimisation considering nonlinear electromechanics, SMO in print.

The Design problem

Fabrication principle and the engineering objective I

· Layer-by-layer design layout: engineering principles



- | \mathcal{B}_0 is comprised of the union of N_{L} elastomeric layers, $\{\mathcal{B}_0^1, \ldots, \mathcal{B}_0^{N_L}\}.$
- II Electrodes intercalated between adjacent elastomer layers.
- $\begin{array}{ll} \mbox{III} & \mbox{Negligible thickness of electrodes,} \\ & \mbox{placed at surface regions} \\ & \{\mathcal{E}^1, \dots, \mathcal{E}^{N_L-1}\}, \end{array}$

$$\mathcal{E}^{i} = \bar{\mathcal{B}}_{0}^{i} \cap \bar{\mathcal{B}}_{0}^{i+1};$$

$$i = \{1, \dots, N_{\mathsf{L}} - 1\}.$$

IV Electrode regions subjected to same voltage value, either positive or negative.

Design in Electro-Mechanics

The Design problem

Fabrication principle and the engineering objective II

· Layer-by-layer design layout: the objective







The Design problem

The design problem: computational considerations I

· Layer-by-layer design layout: computational considerations







Optimal electrode Layout

Control finite elasticity

The Design problem

Regularisation of the problem II

I Smooth representation of electrodes within \mathcal{E}^i via an in-surface phase-field function $\psi_{\mathcal{E}^i}$.

$$\mathbb{V}^{\psi}\varepsilon^{i} = \left\{\psi_{\mathcal{E}^{i}}: \mathcal{E}^{i} \to [0,1], \qquad \psi_{\mathcal{E}^{i}} \in H^{1}(\mathcal{E}^{i})\right\}; \quad i = \{1, \dots, N_{L-1}\},$$



Optimal electrode Layout

Design in Electro-Mechanics

The Design problem

Regularisation of the problem III

I Smooth representation of electrodes within \mathcal{E}^i via an in-surface phase-field function $\psi_{\mathcal{E}^i}$.

 $\mathbb{V}^{\psi}\mathcal{E}^{i} = \left\{\psi_{\mathcal{E}^{i}}: \mathcal{E}^{i} \rightarrow [0,1], \quad \psi_{\mathcal{E}^{i}} \in H^{1}(\mathcal{E}^{i})\right\}; \quad i = \{1, \dots, N_{L-1}\}$

II Smooth extension of $\psi_{E^{l}}$ into the volume yielding the phase field function $\psi_{B_{0}}$.

$$\begin{split} \Lambda(\psi_{\mathcal{E}}, \psi_{\mathcal{B}_{0}}^{*}) &= \inf_{\psi_{\mathcal{B}_{0}} \in \mathbb{V}^{\psi_{\mathcal{B}_{0}}}} \left\{ \int_{\mathcal{B}_{0}} \frac{1}{2} |\nabla_{0}\psi_{\mathcal{B}_{0}}|^{2} dV \right\} \\ \mathbb{V}^{\psi_{\mathcal{B}_{0}}} &= \begin{cases} \psi_{\mathcal{B}_{0}} : \mathcal{B}_{0} \to [0, 1] , & \psi_{\mathcal{B}_{0}} \in H^{1}(\mathcal{B}_{0}) \\ \text{s.t.} & \psi_{\mathcal{B}_{0}} = \psi_{\mathcal{E}^{i}} \text{ on } \mathcal{E}^{i}, \ i = \{1, \dots, N_{L} - 1\} \end{cases} \end{split}$$



Design in Electro-Mechanics

The Design problem

Regularisation of the problem III

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$$\mathbb{V}^{\psi} \varepsilon^{i} = \left\{ \psi_{\mathcal{E}^{i}} : \mathcal{E}^{i} \to [0,1], \quad \psi_{\mathcal{E}^{i}} \in H^{1}(\mathcal{E}^{i}) \right\}; \quad i = \{1, \dots, N_{L-1}\}$$

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$$\begin{split} \Lambda(\boldsymbol{\psi}_{\mathcal{E}}, \boldsymbol{\psi}_{\mathcal{B}_{0}}^{*}) &= \inf_{\boldsymbol{\psi}_{\mathcal{B}_{0}} \in \mathbb{V}^{\boldsymbol{\psi}_{\mathcal{B}_{0}}}} \left\{ \int_{\mathcal{B}_{0}} \frac{1}{2} |\boldsymbol{\nabla}_{0}\boldsymbol{\psi}_{\mathcal{B}_{0}}|^{2} dV \right\} \\ \mathbb{V}^{\boldsymbol{\psi}_{\mathcal{B}_{0}}} &= \begin{cases} \boldsymbol{\psi}_{\mathcal{B}_{0}} : \mathcal{B}_{0} \to [0, 1], & \boldsymbol{\psi}_{\mathcal{B}_{0}} \in H^{1}(\mathcal{B}_{0}) \\ \text{s.t.} & \boldsymbol{\psi}_{\mathcal{B}_{0}} = \boldsymbol{\psi}_{\mathcal{E}^{i}} \text{on } \mathcal{E}^{i}, \ i = \{1, \dots, N_{L} - 1\} \end{cases} \end{split}$$



The Design problem

Regularisation of the problem IV

III Energy Interpolation Scheme

Additive decomposition:

$$e(\psi_{\mathcal{B}_0}(X,\psi_{\mathcal{E}}),F,D_0)=e(F)+\bar{e}_{em}(\psi_{\mathcal{B}_0}(X,\psi_{\mathcal{E}}),F,D_0)$$

· Conditions to be satisfied:

$$\begin{split} \bar{e}_{em}(\psi_{\mathcal{B}_0}(\boldsymbol{X}, \boldsymbol{\psi}_{\mathcal{E}}), \boldsymbol{F}, \boldsymbol{D}_0) \big|_{\psi_{\mathcal{B}_0} = 1} &= e_{em}(\boldsymbol{F}, \boldsymbol{D}_0) \\ \bar{e}_{em}(\psi_{\mathcal{B}_0}(\boldsymbol{X}, \boldsymbol{\psi}_{\mathcal{E}}), \boldsymbol{F}, \boldsymbol{D}_0) \big|_{\psi_{\mathcal{B}_0} = 0} &= \frac{1}{\alpha} e_{em}(\boldsymbol{F}, \boldsymbol{D}_0) \end{split}$$



Proposed scheme

$$\bar{e}_{em}(\psi_{\mathcal{B}_0}(\boldsymbol{X}, \boldsymbol{\psi}_{\mathcal{E}}), \boldsymbol{F}, \boldsymbol{D}_0) = \frac{1}{\left(\psi_{\mathcal{B}_0}(\boldsymbol{X}, \boldsymbol{\psi}_{\mathcal{E}})\right)^p + \alpha(1 - \left(\psi_{\mathcal{B}_0}(\boldsymbol{X}, \boldsymbol{\psi}_{\mathcal{E}})\right)^p)} e_{em}(\boldsymbol{F}, \boldsymbol{D}_0)$$

The Design problem

The regularised problem

The PDE and boundary conditions



The Design problem

The optimal design problem I



$$\mathcal{L}(\psi_{\mathcal{B}_0}(\boldsymbol{\psi}_{\mathcal{E}}), \boldsymbol{\phi}, \varphi, \boldsymbol{p}_{\boldsymbol{\phi}}, p_{\varphi}) = \mathcal{J}(\boldsymbol{\phi}) - \mathcal{W}^{\mathsf{mec}}(\psi_{\mathcal{B}_0}(\boldsymbol{\psi}_{\mathcal{E}}), \boldsymbol{\phi}, \varphi, \boldsymbol{p}_{\boldsymbol{\phi}}) - \mathcal{W}^{\mathsf{ele}}(\psi_{\mathcal{B}_0}(\boldsymbol{\psi}_{\mathcal{E}}), \boldsymbol{\phi}, \varphi, p_{\boldsymbol{\varphi}})$$

$$\mathcal{W}^{\mathsf{rec}}(\psi_{\mathcal{B}_{0}}(\psi_{\mathcal{E}}),\phi,\varphi,p_{\phi}) = \int_{\Omega_{0}} P(\psi_{\mathcal{B}_{0}},F,D_{0}): \nabla_{0}p_{\phi} \, dV - \left(\int_{\Omega_{0}} f_{0} \cdot p_{\phi} \, dV + \int_{\partial\Omega_{0Nm}} t_{0} \cdot p_{\phi} \, dA\right)$$
$$\mathcal{W}^{\mathsf{elec}}(\psi_{\mathcal{B}_{0}}(\psi_{\mathcal{E}}),\phi,\varphi,p_{\phi}) = \int_{\Omega_{0}} D_{0}(\psi_{\mathcal{B}_{0}},F,E_{0}) \cdot \nabla_{0}p_{\varphi} \, dV + \left(\int_{\Omega_{0}} \rho_{0}p_{\varphi} \, dV + \int_{\partial\Omega_{0Ne}} \omega_{0}p_{\varphi} \, dA\right)$$

Sensitivity of the Lagrangian

$$D\mathcal{L}[\Delta\psi_{\mathcal{E}^i}] = \int_{\mathcal{E}^i} g_{\mathcal{E}^i} \, dA$$

Seminar: Control in Times of Crisis

The Design problem

The optimal design problem II

The Optimal Design problem

$$\begin{split} \min_{\psi_{\mathcal{E}^{i}}(X)} \ \mathcal{J}(\boldsymbol{\phi}) \\ \text{s.t} \begin{cases} \text{Statements } \mathcal{W}^{\text{mec}} = 0 \ \mathcal{W}^{\text{elec}} = 0 \\ \text{Constitutive model and energy} \\ \text{interpolation scheme} \\ \text{Laplacian extension into the volume} \\ 0 \leq \psi_{\mathcal{E}^{i}}(X) \leq 1; \ \ \psi_{\mathcal{E}^{i}}(X) \in H^{1}(\mathcal{E}^{i}) \end{cases} \end{split}$$



Phase-field functions

Optimal electrode Layout

$$\begin{split} \partial_{\tau}\psi_{\mathcal{E}^{i}} &= \kappa \boldsymbol{\nabla}_{\mathcal{E}_{i}}^{2}\psi_{\mathcal{E}_{i}} - \partial_{\psi_{\mathcal{E}^{i}}}\left(\frac{1}{4}\Phi(\psi_{\mathcal{E}^{i}}) + \eta \, g_{\mathcal{E}^{i}}(\boldsymbol{\xi})h(\psi_{\mathcal{E}^{i}})\right) \\ 0 &= \boldsymbol{\nabla}_{\mathcal{E}^{i}}\psi_{\mathcal{E}^{i}} \cdot N|_{\partial\mathcal{E}^{i}} \\ \psi_{\mathcal{E}^{i}}|_{\tau_{0}} &= \psi_{\mathcal{E}^{i}}^{0} \\ \Phi(\psi_{\mathcal{E}^{i}}) &= \psi_{\mathcal{E}^{i}}^{2}(1 - \psi_{\mathcal{E}^{i}})^{2} \\ h(\psi_{\mathcal{E}^{i}}) &= \psi_{\mathcal{E}^{i}}^{3}(6\psi_{\mathcal{E}^{i}}^{2} - 15\psi_{\mathcal{E}^{i}} + 10) \end{split}$$



Numerical examples

Numerical example I





J: Martínez-Frutos, R. Ortigosa and A. J: Gil In-silico design of electrode meso-arquitecture for shape morphing dielectric elastomers, IJMPS, submitted.

Numerical examples

Numerical example II



J: Martínez-Frutos, R. Ortigosa and A. J: Gil In-silico design of electrode meso-arquitecture for shape morphing dielectric elastomers, IJMPS, submitted.