Local null controllability of a fluid-rigid | body interaction problem with Navier slip boundary conditions

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FSI problems

Fluid structure problems are coupling systems that involve generally a fluid and rigid or deformable structures.

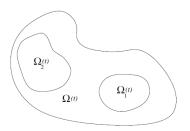


Figure - Fluid-solid interaction problem

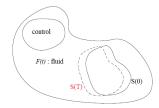
Objectives

The system writes:

$$\begin{cases} z' = Az + Bv + F(z), \\ a'(t) = Cz(t), \\ z(0) = z^{0}, \\ a(0) = a^{0}. \end{cases}$$

- z stands for the velocity of the fluid and the structure velocities and
 a stands for the position of the rigid body.
- v is the control.

Objectives: Find a distributed control v such that: z(T) = 0 and $a(T) = a_T$.



Specificities of this work

- FSI problems : free boundary problem.
- Navier boundary conditions.
- Non linear problem.
- Coupled system.

Boundary conditions

The non-slip (or Dirichlet) boundary conditions :

$$U_f = \mathbf{v_s}$$
 on $\partial \Omega$.

The slip (or Navier) boundary conditions :

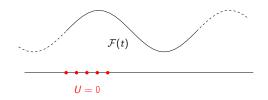
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 \begin{cases} (U_f - \mathbf{v_s})_n = 0 & \text{on } \partial \Omega, & \text{(Impermeability condition)} \\ [\mathbb{T}(U_f, P_f) \mathbf{n} + \beta (U_f - \mathbf{v_s})]_{\tau} = 0 & \text{on } \partial \Omega, & \text{(Slip condition)} \end{cases}
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- $\mathbb{T}(U_f, P_f)n$: The force exerted by the fluid on the surface.
- \bullet β : The friction coefficient.

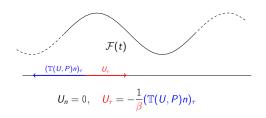
In the FSI problems, the vector U_f designates the fluid velocity and the vector v_s stands for the structure velocity.

Boundary conditions

Dirichlet boundary conditions



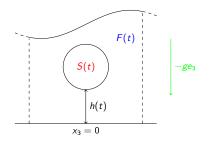
Navier boundary conditions



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The influence of the boundary conditions

Rigid structure



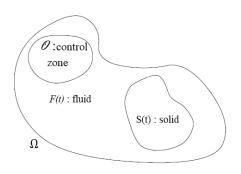
- M. Hillairet, D. Gérard-Varet, C. Wang. 2014 :
 - Dirichlet boundary conditions : No contact at x = 0.
 - Navier boundary conditions : Collision can occur at x = 0.

Previous works with Navier boundary conditions

Existence results:

- Existence of weak solutions : D. Gérard-Varet, M. Hillairet. 2014
- Uniqueness of the weak solution: N V. Chemetov, Š. Nečasová, B. Muha. 2017.
- Existence and uniqueness of strong solution: C. Wang. 2014

Problem setting



- ullet $\mathcal{F}(t)$: the fluid domain.
- S(t): the rigid body.
- $\bullet \ \Omega = \mathcal{F}(t) \cup S(t).$

$$\partial \mathcal{F}(t) = \partial S(t) \cup \partial \Omega$$

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Fluid equations

The incompressible Navier-Stokes system

$$\begin{cases} \partial_t U + (U \cdot \nabla)U - \nabla \cdot \mathbb{T}(U, P) = v \mathbf{1}_{\mathcal{O}}, & \text{in } (0, T) \times \mathcal{F}(t), \\ \nabla \cdot U = 0, & \text{in } (0, T) \times \mathcal{F}(t), \end{cases}$$

- $\mathcal{F}(t) \subset \mathbb{R}^3$: Fluid domain that depends on time.
- $U(t,x) \in \mathbb{R}^3$: The fluid velocity.
- P(t,x): The fluid pressure.

The non linear term :

$$(U \cdot \nabla)U = \sum_{i=1}^{3} U_{i} \frac{\partial U}{\partial x_{i}}.$$

The Cauchy stress tensor

$$\mathbb{T}(U,P) = -PI_2 + 2\nu D(U),$$

u: The viscosity of the fluid.

$$[D(U)]_{i,j} = \frac{1}{2} \left(\frac{\partial U_i}{\partial x_i} + \frac{\partial U_j}{\partial x_i} \right).$$

Structure equations

The structure position:

$$X_{\mathcal{S}}(t,y) = h(t) + R_{\theta(t)}y, \quad y \in \mathcal{S}.$$

h(t): The solid position of the solid.

 θ : The rotation of the solid.

The structure velocity

$$U_S(t,x) = h'(t) + \omega(t)(x - h(t))^{\perp}, \quad x \in \partial S(t),$$

 $\omega(t)$: The angular velocity of the solid.

Newton's laws:

$$\left\{ \begin{array}{c} mh''(t) = -\int_{\partial S(t)} \mathbb{T}(U,P) n \ d\Gamma & t \in (0,T), \\ J\omega'(t) = -\int_{\partial S(t)} (x-h(t))^{\perp} \cdot \mathbb{T}(U,P) n \ d\Gamma & t \in (0,T). \end{array} \right.$$

m > 0 is the mass of the rigid structure and J > 0 its moment of inertia

Navier boundary conditions

On the fixed boundary part $\partial\Omega$, the boundary conditions write

$$\left\{ \begin{array}{ll} U_n = 0 & \text{ on } (0,T) \times \partial \Omega,, \\ \left[2 \nu D(U) n + \beta_\Omega U \right]_\tau = 0 & \text{ on } (0,T) \times \partial \Omega, \end{array} \right.$$

Navier boundary conditions

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On the moving boundary $\partial S(t)$, the boundary conditions write

$$\left\{ \begin{array}{ll} (U - \textcolor{red}{U_S})_n = 0 & \text{ on } (0, T) \times \partial S(t), \\ [2\nu D(U)n + \beta_S (U - \textcolor{red}{U_S})]_{\tau} = 0 & \text{ on } (0, T) \times \partial S(t). \end{array} \right.$$

where β_{Ω} and β_{S} are the friction coefficients.

Navier boundary conditions

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where β_{Ω} and $\beta_{\mathcal{S}}$ are the friction coefficients.

Initial conditions

$$U(0,x) = u^{0}(x), x \in \mathcal{F}(0), \quad h(0) = h^{0}, \quad h'(0) = h^{1}, \quad \omega(0) = \omega^{0},$$

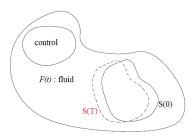
Problem setting

$$\begin{cases} \partial_t U + (U \cdot \nabla)U - \nabla \cdot \mathbb{T}(U, P) = v \mathbf{1}_{\mathcal{O}} & t \in (0, T), \ x \in \mathcal{F}(t), \\ \nabla \cdot U = 0 & t \in (0, T), \ x \in \mathcal{F}(t), \end{cases} \\ \begin{cases} mh''(t) = -\int_{\partial S(t)} \mathbb{T}(U, P) n \ d\Gamma & t \in (0, T), \\ J\omega'(t) = -\int_{\partial S(t)} (x - h(t))^{\perp} \cdot \mathbb{T}(U, P) n \ d\Gamma & t \in (0, T), \end{cases} \\ \begin{cases} U_n = 0 & t \in (0, T), \ x \in \partial \Omega, \\ [2\nu D(U)n + \beta_{\Omega}U]_{\tau} = 0 & t \in (0, T), \ x \in \partial S(t), \\ [2\nu D(U)n + \beta_{S}(U - U_{S})]_{\tau} = 0 & t \in (0, T), \ x \in \partial S(t), \end{cases} \end{cases}$$

$$U(0,x) = u^{0}(x), x \in \mathcal{F}(0), \quad h(0) = h^{0}, \quad h'(0) = h^{1}, \quad \omega(0) = \omega^{0},$$

Problem setting

Objective: Find a control v such that the velocities are equal to zero at final time T.



Main result

Theorem

Assume that $\beta_S > 0$. There exists $v \in L^2(0, T; [L^2(\mathcal{O})]^3)$ such that

$$U(T,.)=0$$
 in $\mathcal{F}(T), \quad h(T)=0, \quad h'(T)=0$
$$\omega(T)=0, \quad \theta(T)=0,$$

provided that the velocities are small enough and the final state is close enough to the initial state.

Note: If $h(T) = h_T$ and $\theta(T) = \theta_T$, we can reduce the problem to h(T) = 0, $\theta(T) = 0$.

Previous null controllability results

Dirichlet boundary conditions:

- O. Imanuvilov, T. Takahashi. 2007.
- M. Boulakia and A. Osses. 2007.
- M. Boulakia and S. Guerrero. 2011.

Outline

Change of variable

2 Linearized system

Fixed-point

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Change of variable

2 Linearized system

Fixed-point

Change of variables

T. Takahashi. 2003.

• Change of variables : $X(t,\cdot):\Omega\longrightarrow\Omega$ such that

$$X(t,\cdot) = X_{S}(t,\cdot) : S = S(T) \longrightarrow S(t)$$

$$Y(t,\cdot) = X^{-1}(t,\cdot)$$

$$u(t,\cdot) = \nabla Y(t,X(t,\cdot))U(t,X(t,\cdot)), \quad p(t,\cdot) = P(t,X(t,\cdot)).$$

Change of variables

C. Wang. 2014.

Fluid equations:

$$u(t,y) = (\nabla Y)(t,X(t,y))U(t,X(t,y)), \quad p(y,t) = P(t,X(t,y)),$$

$$\begin{cases} \partial_t u - \nabla \cdot \mathbb{T}(u,p) = v1_{\mathcal{O}} + F(u,p) & t \in (0,T), \ y \in \mathcal{F}, \\ \nabla \cdot u = 0 & t \in (0,T), \ y \in \mathcal{F}, \end{cases}$$

Structure equations

$$\begin{cases} mh''(t) = -\int_{\partial S} \mathbb{T}(u, p) n \ d\Gamma & t \in (0, T), \\ J\omega'(t) = -\int_{\partial S} y^{\perp} \cdot \mathbb{T}(u, p) n \ d\Gamma & t \in (0, T), \end{cases}$$

 $u_S = h'(t) + \omega(t) y^{\perp}, \quad y \in \partial S,$

Boundary conditions

$$\begin{cases} u_{n} = 0 & t \in (0, T), y \in \partial \Omega, \\ \left[2\nu D(u)n + \beta_{\Omega}u\right]_{\tau} = 0 & t \in (0, T), y \in \partial \Omega, \\ (u - u_{S})_{n} = 0 & t \in (0, T), y \in \partial S, \\ \left[2\nu D(u)n + \beta_{S}(u - u_{S})\right]_{\tau} = 0 & t \in (0, T), y \in \partial S. \end{cases}$$

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Linearized system

$$\begin{cases} \partial_t u - \nabla \cdot \mathbb{T}(u, p) = f + v \mathbf{1}_{\mathcal{O}} & t \in (0, T), \ y \in \mathcal{F}, \\ \nabla \cdot u = 0 & t \in (0, T), \ y \in \mathcal{F}. \end{cases}$$

$$\begin{cases} mh''(t) = -\int_{\partial S} \mathbb{T}(u, p) n \ d\Gamma & t \in (0, T), \\ J\omega'(t) = -\int_{\partial S} y^{\perp} \cdot \mathbb{T}(u, p) n \ d\Gamma & t \in (0, T). \end{cases}$$

$$u_S = h'(t) + \omega(t) y^{\perp}, \quad y \in \partial S.$$

Boundary conditions

$$\begin{cases} u_{n} = 0 & t \in (0, T), \ y \in \partial \Omega, \\ \left[2\nu D(u)n + \beta_{\Omega}u\right]_{\tau} = 0 & t \in (0, T), \ y \in \partial \Omega, \\ (u - u_{S})_{n} = 0 & t \in (0, T), \ y \in \partial S, \\ \left[2\nu D(u)n + \beta_{S}(u - u_{S})\right]_{\tau} = 0 & t \in (0, T), \ y \in \partial S. \end{cases}$$

$$u(0, y) = u^{0}(y), \ y \in \mathcal{F}, \quad h(0) = h^{0}, \quad h'(0) = h^{1}, \quad \omega(0) = \omega^{0},$$

Linearized system

The linearized system writes:

$$\begin{cases} u' = Au + Bv + F, \\ a'(t) = Cu(t), \\ u(0) = u^{0}, \\ a(0) = a^{0}. \end{cases}$$

A is introduced in C. Wang. 2014,

$$a = (h, \theta), \quad Cu = (h', \omega).$$

Theorem (Null controllability of the linearized system)

For all F, u^0 , a^0 there exists a control $v \in L^2(0, T; [L^2(\mathcal{O})]^2)$ such that u(T) = 0 and a(T) = 0.

Main steps of the proof

We use the theory of O. Imanuvilov, T. Takahashi. 2007.

- Main ingredient : Carleman estimate for the adjoint system of the linearized problem.
- Difficulties :
 - Control the velocity of the rigid body with the fluid velocity:

$$||u_S \cdot n||_{L^2(\partial S)} \leq C ||u||_{L^2(\mathcal{F})},$$

$$\beta_{S} \|u_{S} \cdot \tau\|_{L^{2}(\partial S)} \leq C \left(\|u\|_{L^{2}(\partial \mathcal{F})} + \|u\|_{L^{2}(\mathcal{F})} + \|(\nabla \times u)_{\tau}\|_{L^{2}(\partial \mathcal{F})} \right).$$

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Fixed-point

$$\begin{cases} \partial_{t}u - \nabla \cdot \mathbb{T}(u, p) = v1_{\mathcal{O}} + F(u, p) & t \in (0, T), \ y \in \mathcal{F}, \\ \nabla \cdot u = 0 & t \in (0, T), \ y \in \mathcal{F}, \end{cases}$$

$$u_{S} = h'(t) + \omega(t)y^{\perp}, \quad y \in \partial S,$$

$$\begin{cases} mh''(t) = -\int_{\partial S} \mathbb{T}(u, p)n \ d\Gamma & t \in (0, T), \\ J\omega'(t) = -\int_{\partial S} y^{\perp} \mathbb{T}(u, p)n \ d\Gamma & t \in (0, T), \end{cases}$$

$$\begin{cases} u_{n} = 0 & t \in (0, T), \ y \in \partial \Omega, \\ [2\nu D(u)n + \beta_{\Omega}u]_{\tau} = 0 & t \in (0, T), \ y \in \partial S, \\ [2\nu D(u)n + \beta_{S}(u - u_{S})]_{\tau} = 0 & t \in (0, T), \ y \in \partial S. \end{cases}$$

There exists a function $\rho \in C([0, T])$ such that the application

$$\Phi: f \longrightarrow F(u, p)$$

defines a contraction on

$$K = \left\{ f \in L^2(0,T;[L^2(\mathcal{F})]^2), \quad \left\| \frac{f}{\rho} \right\|_{L^2(0,T;[L^2(\mathcal{F})]^2)} \le R \right\} \text{ and } \Phi(K) \subset K$$

Conclusion and perspectives

- Global controllability.
- Reduce the controls.

Navier Stokes system: S.Guerrero and C.Montoyay. 2017

• Deformable structure : Dirichlet boundary condition : J.Lequeurre 2013

Thank you for your attention